

# B.A. Taking the $\delta$ -Bin into Account

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## Outline

Motivation

Double Counting,  $\sum_{P \neq 0} \int d^4x$ , Thr. Expn

Pull-up hep-ph/0102257 [predates Lsæt]  
Hoang, I.S., Stewart

$\delta$ -bin

Examples: I  $\delta$ -bin  $\Rightarrow$  pull-up } NRQCD  
II resolves pinch sing. issue } [unphysical]

pinches in soft boxes  
}}, }}

$\int \frac{dk^0}{(k^0 + i\epsilon)(k^0 - i\epsilon)}$   $\leftarrow$  problem in Thr. Expn  
[soft mode can't be defined]

III SCET<sub>I</sub> loops  
IV tree level

V SCET<sub>I</sub> implications <sup>for</sup> endpoint sing. [no unphysical ...]

$$\int_0^1 dx \frac{\delta_{II}(x)}{x^2}$$

pinch surfaces  
collinear horizon

[No change to ~~cases~~ that work.]

- Facts
- EFT is indep. of IR regulator choice [d.o.f.]
  - Poor IR ( $\pm UV$ ) regulators can spoil p.c. [not an issue]
  - $\int \frac{d^4 k}{k^4} = \frac{i}{16\pi^2} \left( \frac{1}{\epsilon_{UV}} - \frac{1}{\epsilon_{IR}} \right)$   
 $\uparrow$  requires  $\delta Z$

Motivation

- Formulate NRCQD, SCET indep. of dim. reg. (UV) and IR regulator
- Solve dble counting, pull-up, singularity issues

- Lamb Shifts,  $e\bar{e}$  production
- Static Scattering
- $B \rightarrow \pi e \bar{e}$
- $B \rightarrow \pi\pi$  (annih.)
- $\gamma^* \pi \rightarrow \pi$  (subl.)
- $\gamma^* p \rightarrow p$  ( $F_2$ )
- $\gamma^* \rightarrow q \bar{q} g$
- $B^0 \rightarrow D^0 K^0$
- 
-

- Insert picture from (4) -

(3)

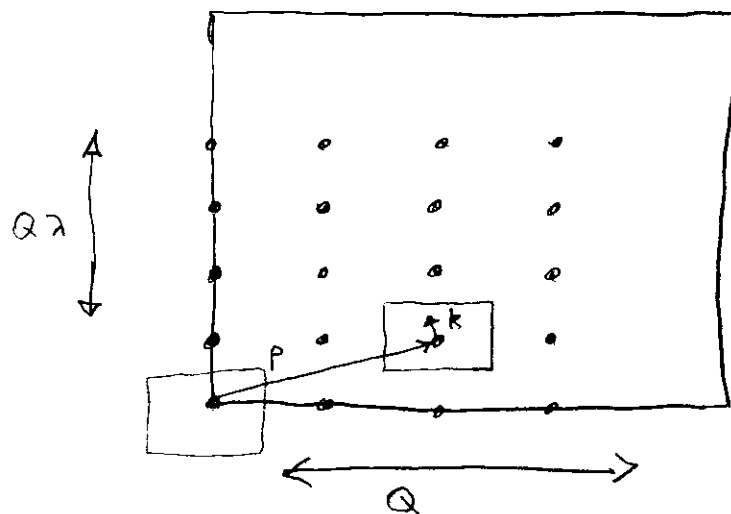
Why two low energy d.o.f.? p.c.

Must capture all IR scales  $m_1, m_2, \dots$

[ Sometimes a d.o.f makes p.c. so opaque that it's impossible to setup the theory ]

$$\psi(x) = \sum_p e^{-ip \cdot x} \psi_p(x)$$

collinear  $P = p + p_r$   
 label residual

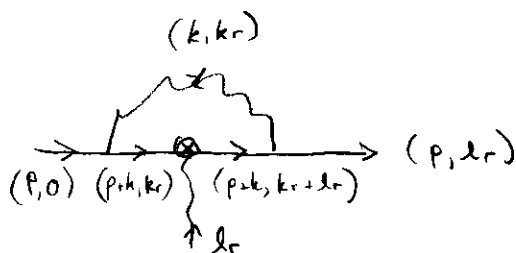


- manifest p.c.
- homo. derivatives

Note:  $\sum_{p \neq 0}$  otherwise  $P$  is soft!

this is the zero-bin

Loop

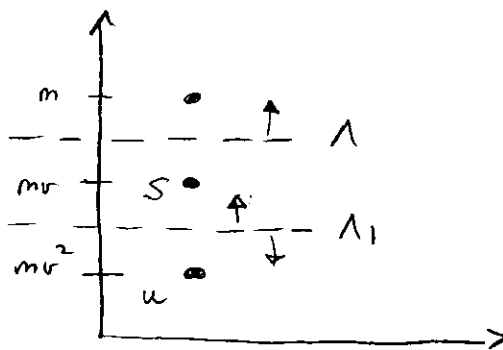
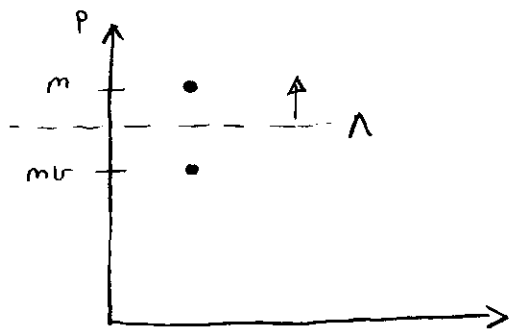


$$\sum_k \int dkr = \int dk \quad [LMR]$$

leave  $lr$

[ later, but formulation at level of int. is incorrect. ]

Double Counting



cutoffs

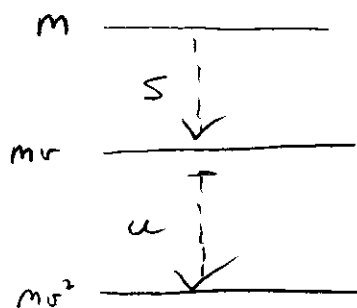
Scaleless Reg.

[double counting in UV taken care of by W. Coeff.]

$\Lambda_1 \rightarrow 0$  for S double counts IR with UV of u

$$I_s = \frac{A}{E_{uv}} + \frac{B}{E_{IR}} + f\left(\frac{s}{\mu}\right)$$

$$I_u = \frac{-B}{E_{uv}} + \frac{C}{E_{IR}} + g\left(\frac{u}{\mu}\right)$$



one mode at a time

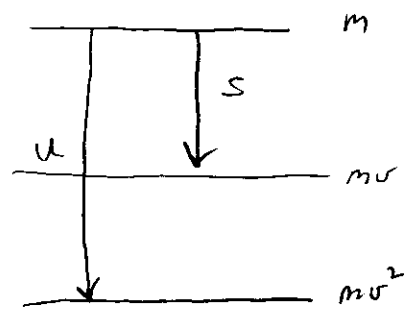
wrong with both modes

- doesn't reproduce IR
- doesn't run simult. from UV [modes not formulated simult.]

[GFT, cot.]

Instead

Adds  $B \left( \frac{1}{E_{uv}} - \frac{1}{E_{IR}} \right)$  to S



subtract from u

$$I_s = \frac{A+B}{E_{uv}} + \dots$$

$$I_u = \frac{-B}{E_{uv}} + \frac{C}{E_{IR}}$$

by hand

pulled up to hard scale

Strong Check:

NLL	NRQED	Anom. Dim	
$d^3 \ln^3 d$	Landr		H, e <sup>+</sup> e <sup>-</sup>
$d^7 \ln^2 d$	Landr, h.f.s		"
$d^3 \ln^2 d$	$\Delta \Gamma$		e <sup>+</sup> e <sup>-</sup>

$\phi$ -Bin Subtraction:

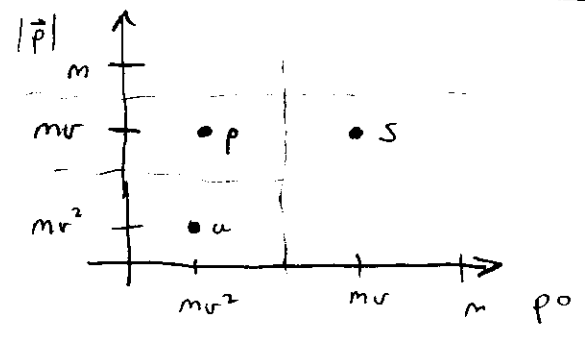
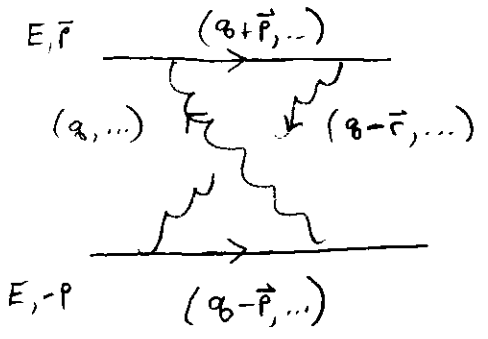
unique

$$\sum_{\{p_i \neq 0\}} \int \prod_i d^4 k_i F(p_i, k_i) \rightarrow \int \prod_i d^4 p_i \left[ F(p_i) - \sum_{j \in U} F_j^{sub}(p_i) \right]$$

$\uparrow$   
 labelled  
 determined by field contractions  
 - but with mom. conservation it  
 acts on entire loop integral

$\uparrow$   
 removes support  
 of integrand  
 to avoid dble.  
 counting!

NRQCD



[ soft - usoft gluons  
 ptal - soft quarks

$$I_s = \sum_{\substack{q^0 \neq 0 \\ q^4 \neq 0 \\ q^4 \neq \vec{r}}} \int d^4 q_r \frac{1}{(q^0 + ic)(q^0 - ic)(q^2)(q - \vec{r})^2}$$

$$I_s = \tilde{I}_s - I_1 - I_2$$

$\uparrow$   
 naive

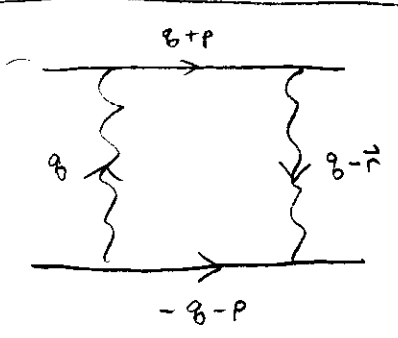
$$\tilde{I}_S = \int \frac{d^d q}{P_{en}}$$

$$I_1 = \int \frac{d^d q}{(q^0)(q^0)(q^2)(-\vec{r}^2)} \quad \leftarrow \quad q^0 = 0$$

$$I_2 = \int \frac{d^d q}{(q^0)(q^0)(-\vec{r}^2)(q-\vec{r})^2} \quad \leftarrow \quad q^0 = \vec{r}$$

$$\tilde{I}_S = \frac{1}{E_{IR}} + \ln\left(\frac{\mu^2}{\vec{r}^2}\right)$$

$$-I_1 - I_2 = -\frac{1}{E_{IR}} + \frac{1}{E_{UV}} \quad \checkmark \quad \text{pull-up (I)}$$



C-box arose ↓

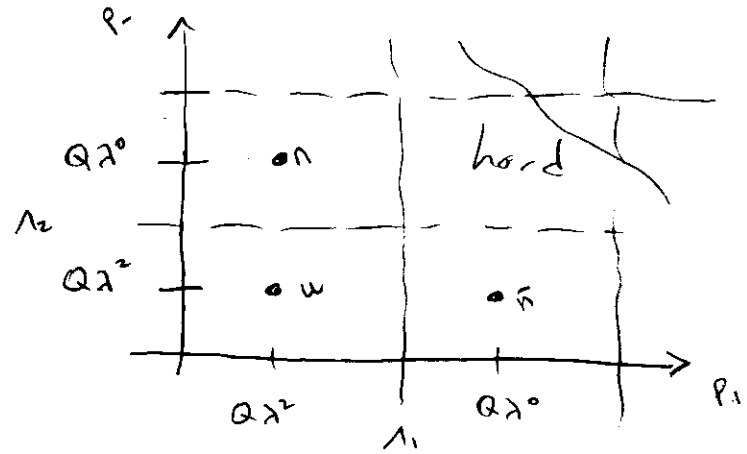
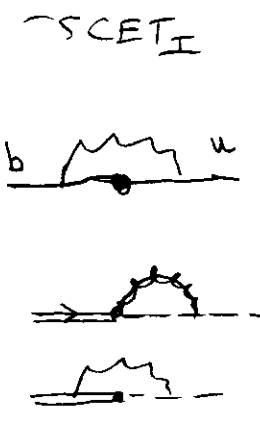
$$I_S = \sum_{\substack{q^0 \neq 0 \\ q^0 \neq 0 \\ q^0 \neq \vec{r}}} \int d^d q \frac{1}{(q^0+i\epsilon)(-q^0+i\epsilon)(q^2)(q-\vec{r})^2}$$

$$I_S = \underbrace{\tilde{I}_S - I_1 - I_2}_{\text{like before}} - \underbrace{I_3}_{p^0=0} + \underbrace{I_{13} + I_{23}}_{p^0=0 \text{ limit of } I_1, I_2}$$

$$I_3 = \int \frac{d^d q}{(q^0+i\epsilon)(-q^0+i\epsilon)(-\vec{q}^2)(-(q-\vec{r})^2)}$$

$\left. \begin{array}{l} \tilde{I}_S - I_3 \\ I_1 - I_{13} \\ I_2 - I_{23} \end{array} \right\}$  have no pinch!

[ Equivalent to "ignore the pinch" etc. ]



Check IR [diff. regulators]

n-collin vs. usoft

$$\begin{aligned}
 I_{\text{full}} &= \int \frac{d^4 q}{(q^2) (q^2 + 2p_+ \cdot q + i\epsilon) (q^2 + 2p_- \cdot q)} \quad \text{on-shell} \\
 &= \int_{-p^-}^0 d q^- \int d^2 q_{\perp} \frac{q^-}{(q_{\perp}^2 + q^{-2}) (q_{\perp}^2)} + \text{finite} \\
 &= \text{Li}_2\left(\frac{-R_{\perp}^2}{R_{\perp}^2}\right) + \ln\left(\frac{R_-}{p^-}\right) \ln\left(\frac{q-p_-}{-R_{\perp}^2}\right) + \dots
 \end{aligned}$$

$$\begin{aligned}
 \Lambda_{\perp}^2 &\geq q_{\perp}^2 \geq R_{\perp}^2 \\
 \Lambda^2 &\geq q^2 \geq R^2
 \end{aligned}$$

$$I_c = \sum_{\tilde{q} \neq 0} \int d^4 q \frac{2 \tilde{n} \cdot (p + q)}{(\tilde{n} \cdot q) (q^2 + 2p \cdot q) (q^2)}$$

$$\hat{I}_c = \int d^4 q (\dots) = \int_{-p^-}^0 d q^- \int d^2 q_{\perp} \frac{2 p^-}{(q^-) (p^-) (q_{\perp}^2)} \quad \begin{matrix} q^- = 0 \\ \text{IR div.} \end{matrix}$$

$$I_1 = \int \frac{d^4 q}{(\tilde{n} \cdot q) (\tilde{n} \cdot p + q) (q^2)} = \int_{-\infty}^0 (\dots)$$

$$\tilde{I}_c - I_1 = \int_{-p^-}^{-\infty} (\dots) \quad \text{no } q^- = 0 \text{ div.}$$

indep. of cutoffs

$$- \widehat{I}_c = - \ln \left( \frac{\Omega_{\perp}^2}{\Lambda_{\perp}^2} \right) \ln \left( \frac{\Omega_{-}}{p^{-}} \right) + \dots$$

$$- I_1 = + \ln \left( \frac{\Omega_{\perp}^2}{\Lambda_{\perp}^2} \right) \ln \left( \frac{\Omega_{-}}{\Lambda_{-}} \right) + \dots$$

$$I_{us} = \text{Liz} \left( \frac{-\Omega_{\perp}^2}{\Omega_{-}^2} \right) + \ln \left( \frac{\Omega_{-}}{\Lambda_{-}} \right) \ln \left( \frac{\Omega_{-} \Lambda_{-}}{\Omega_{\perp}^2} \right) + \dots$$

$$\text{sum} = \underbrace{\text{Liz}(\dots) + \ln \left( \frac{\Omega_{-}}{p^{-}} \right) \ln \left( \frac{\Omega_{-} \Lambda_{-}}{\Omega_{\perp}^2} \right)}_{\text{IR agrees}} + \underbrace{\ln^2 \left( \frac{\Lambda_{-}}{p^{-}} \right) - \ln^2 \left( \frac{\Lambda_{-}}{\Lambda_{-}} \right)}_{\text{UV}}$$

IR agrees



need  $I_1$  else

$\ln \left( \frac{\Omega_{\perp}^2}{\Lambda_{\perp}^2} \right) \ln \Omega_{-}$  spoils it

[ Also needed so no IR x UV cross term for c.t. ]

Offshellness  $p^+ \neq 0$

$$I_{\text{full}} = \ln^2 \left( \frac{-p^+}{p^{-}} \right)$$

$$I_{us} = \left( \frac{1}{\epsilon_{uv}^2} + \frac{2}{\epsilon_{uv}} \ln \left( \frac{\mu}{-p^+} \right) + 2 \ln^2 \left( \frac{\mu}{-p^+} \right) \right)$$

$$\widehat{I}_c = \left( \frac{-2}{\epsilon_{uv} \epsilon_{IR}} - \frac{2}{\epsilon_{IR}} \ln \left( \frac{\mu}{-p^2} \right) - \ln^2 \left( \frac{\mu^2}{-p^2} \right) + \left( \frac{2}{\epsilon_{IR}} - \frac{2}{\epsilon_{uv}} \right) \ln \left( \frac{\mu}{p^{-}} \right) \right)$$

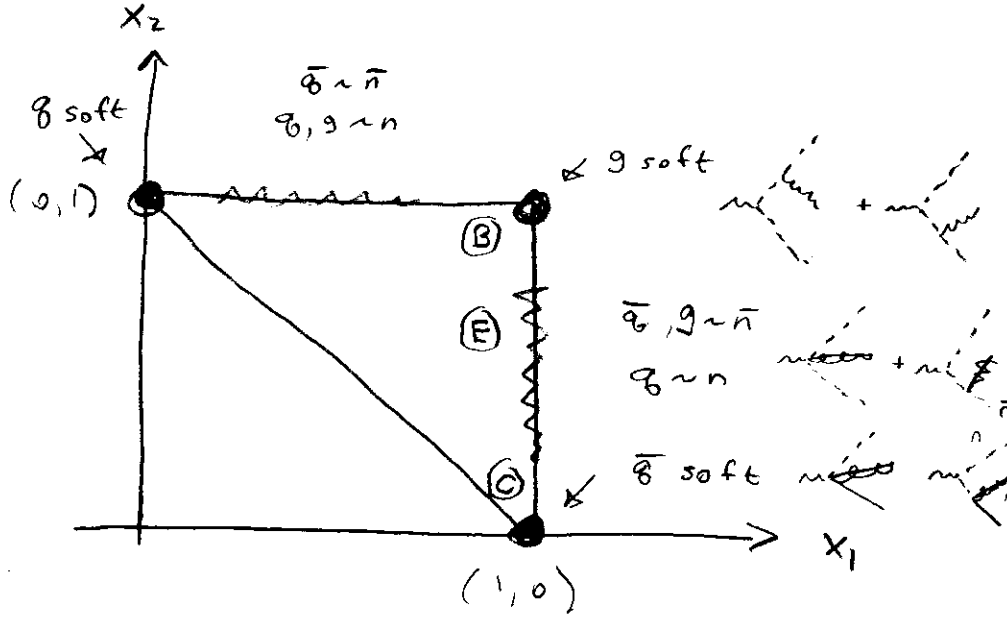
$$- I_1 = \left( \left( \frac{2}{\epsilon_{IR}} - \frac{2}{\epsilon_{uv}} \right) \left( \frac{1}{\epsilon_{uv}} + \ln \left( \frac{\mu^2}{-p^2} \right) - \ln \left( \frac{\mu}{p^{-}} \right) \right) \right)$$

$$\widehat{-EFT} = \frac{-1}{\epsilon_{uv}^2} - \frac{2}{\epsilon_{uv}} \ln \left( \frac{\mu}{p^{-}} \right) - 2 \ln^2 \left( \frac{\mu}{p^{-}} \right) + \ln^2 \left( \frac{-p^+}{p^{-}} \right)$$

$$\gamma^* \rightarrow \bar{q} \bar{q} q$$

1 2 3

$$X_i = \frac{2E_i}{Q}$$



full cross section

$$\frac{d\sigma}{\sigma_0 dx_1 dx_2} = \frac{2 ds}{3\pi} \left[ \frac{x_1^2 + x_2^2}{(1-x_1)(1-x_2)} \right]$$

(B)  $\frac{2}{(1-x_1)(1-x_2)}$

(E)  $\frac{x_2^2 + 1}{(1-x_2)(1-x_1)}$

(C)  $\frac{1}{(1-x_1)}$

$x_1, x_2 \sim 1$

$x_2 \neq 0, x_2 \neq 1, x_1 \sim 1$

$x_2 \sim 0, x_1 \sim 1$

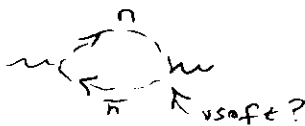
at tree level  
implemented by  
 $\Sigma$  in SCET  
 $p \neq 0$   
no dble counting

↑ p.c. as expected

pick an  $x_1, x_2$ , decide where you are

What if I want to add them?

can't use loops  
ie thr expr.

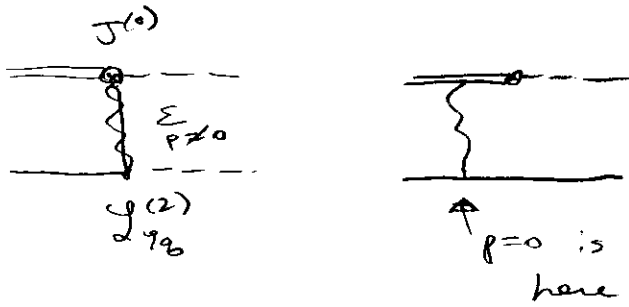


$$\left\{ \frac{\pm k}{n(k+p) \bar{n}(k-p)} = 0 \right.$$

can square m.elt! but must implement subtractions  
in phase space. then sum gives full theory results

- So Subtractions are needed at tree level  
When we integrate over ext. momenta  
Σ are unique

eg.  $B \rightarrow \pi \ell \bar{\nu}$  in SCET<sub>I</sub>



no endpt sing.  
Agrees with QCD

Match  $\Sigma_{p \neq 0}$  onto SCET<sub>II</sub>

Homework: • repeat for SCET<sub>II</sub>

- Fourier transform to get position space formulation of  $\Sigma_{p \neq 0}$